

The Effect of Oblateness on the mean motion of rotating Cyclic Kite Configuration.

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Abstract: This paper deals with the effect of oblateness of the first point mass of the cyclic kite configuration on the mean motion of the rotating frame lying on plane of the kite configuration by using python language .Graphical studies show that the mean motion increases with the increase of the values of oblateness parameter of the primary of mass $m_1 = (1 - \mu) / 2$ where $\mu \in (0, 3^{-1})$. The oblateness of other bodies also may affect the mean motion of the rotating frame.

Keywords: Kite configuration, Cyclic kite configuration, Mass parameter, Rotating frame, Mean motion.

- 1. Definition:** Before to introduce the exact problem; let us define first the kite configuration with classifications. Kite configuration means a quadrilateral with two pairs of equal adjacent sides and cyclic kite configuration is a kite whose vertices lie on the circumference of a common circle. From different articles of previous authors; it is to be noted that a quadrilateral is a kite if and only if their diagonals are perpendicular to each other and at least one of the two diagonals is a line of symmetry. In some cases both the diagonals may be the lines of symmetry.

Thus, the cyclic kite configuration can be classified into two classes.

- i.The kite formed by the combination of one equilateral and one isosceles triangle.
- ii.The kite formed by the combination of two congruent isosceles right angled triangles.

2. Introduction:

The researchers of celestial mechanics used to discuss the particular cases of both the branches such as Three-body, Four-body, Five-body (configuration or problem). In the field of four body configuration MacMillon et al (1932) proved two theorems in detail for the existence of quadrilateral configuration. Brumberg (1957) studied permanent configuration of the Four-body problem.

In (1996) the symmetric central configuration of equal masses was first studied by Albouy. Long et al (2002) proved that a convex non-collinear planar four body central configuration with three equal masses must be a kite .He also analyzed the four-body central configuration with two pairs of equal masses. Chavela et al (2007) invented that if in a convex four body central configuration two equal masses are located at opposite vertices of a quadrilateral and at most the mass of one of the remaining particles is larger than the equal masses then the configuration must be a kite shaped quadrilateral. Further Albouy et al (2008) discussed some properties of a quadrilateral configuration and establish some relations among the masses of the system and it was proved that the planar four-body configuration is a convex central configuration and is symmetric about it's one diagonal if and only if the masses of two particles at the end of the other diagonal are equal. To construct a general four-body problem Pina et al (2009)

developed a new algorithm with the ratios of directed areas and the corresponding scalar areas and in (2010) they developed co-ordinates of four particles of kite configuration in terms of principle moments of inertia and Eulerian angles independently. In Newtonian four-body problem, the cyclic central configuration has been studied for the first time by Cors et al (2012) in which they considered the six mutual distances of the particles as their co-ordinates and showed that the four point masses forming a kite lying on a two-dimensional plane. Also they invented the kite configuration and isosceles trapezoidal configuration. Further they have shown that if any central configuration contains two equal masses then the configuration must have a line of symmetry. Balint Erdi et al (2016) extended the work of Cors et al (2012) in three cases (one in convex case and two in concave cases) by expressing the masses of the central configuration in terms of angle coordinates. Further they claimed that formulas derived represent the exact analytical solutions of the four-body configuration. Deng et al (2017) used mutual distances as the coordinates and proved that the diagonals of a cyclic central quadrilateral cannot be perpendicular except that the configuration is a kite. Corbera et al (2017) proved that the diagonals of a four body central configuration are perpendicular then the configuration must be a kite. Further they verified the same theorem in the four-vertex convex central configuration. Hassan (2023) studied cyclic kite configuration by three theorems and he expressed the masses of the four particles in terms of a mass parameter μ and the total mass M of the system. Also he developed the co-ordinates of the four masses and justified all the results of previous authors by using his results. Not only that, he also derived the mean motion of the rotating frame lying in the plane of the kite. Presently we have proposed to study the effect of the oblateness of the first body of mass $m_1 = \frac{M(1-\mu)}{2}$ on the mean motion of the rotating frame.

3: Theorem:- When the cyclic kite configuration is formed by one equilateral triangle and one isosceles triangle with the common base P_2P_4 :

Statement:-The necessary and sufficient conditions for the existence of cyclic kite configuration with four positive point masses m_k ($k = 1, 2, 3, 4$) at the vertices P_k of the kite are that there exists a non-dimensional mass parameter

$$\mu \in (0, 3^{-1}) \text{ such that } \frac{m_2}{M} = \frac{m_4}{M} = \mu, \quad \text{then} \quad m_1 = \frac{M(1-\mu)}{2}, m_3 = \frac{M(1-3\mu)}{2}, \text{ where} \\ m_1 + m_2 + m_3 + m_4 = M.$$

Proof.

Let P_k ($k = 1, 2, 3, 4$) be the positions (ref. fig 1) of the four positive point masses m_k moving on a common circular orbit with center at O, radius R and diameter $P_1OP_3 = 2R$ as the axis of symmetry. Considering the center of the circle as the origin, the axis of symmetry as the x-axis and a line YOY' perpendicular to P_1OP_3 as the y-axis. Using the properties of cyclic quadrilateral $P_1P_2P_3P_4$ in the form of cyclic kite configuration and the cyclic equilateral triangle $P_1P_2P_4$, the coordinates of the four point masses can be written as $P_k(R \cos \theta_k, R \sin \theta_k)$, where $\theta_1 = 0^\circ, \theta_2 = 120^\circ, \theta_3 = 180^\circ, \theta_4 = 240^\circ$.

Thus the position vectors of the four point masses are given by $\vec{r}_k = R \cos \theta_k \hat{i} + R \sin \theta_k \hat{j}$

$$\Rightarrow \vec{r}_1 = R\hat{i}, \vec{r}_2 = -\frac{R}{2}\hat{i} + \frac{\sqrt{3}}{2}R\hat{j}, \vec{r}_3 = -R\hat{i}, \vec{r}_4 = -\frac{R}{2}\hat{i} - \frac{\sqrt{3}}{2}R\hat{j}. \quad (1)$$

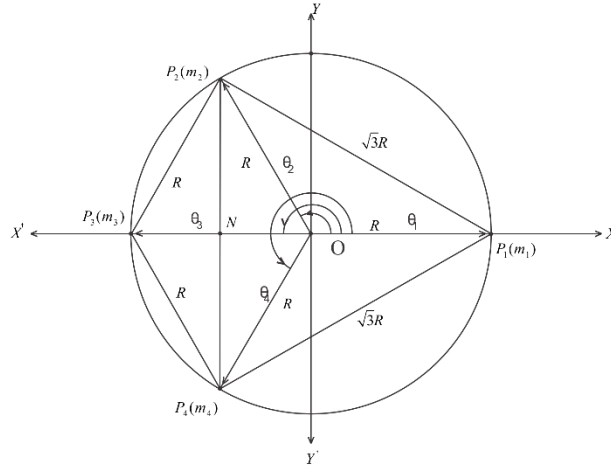


Figure-1

The position vector of the center of mass of the system of four- point masses is given by

$$\vec{c} = M^{-1} \sum_{k=1}^4 m_k \vec{r}_k. \quad (2)$$

If we consider the center' of mass as the origin, then $\vec{c} = \hat{0} \Rightarrow m_1 \vec{r}_1 + m_2 \vec{r}_2 + m_3 \vec{r}_3 + m_4 \vec{r}_4 = \hat{0}$

$$\begin{aligned} \Rightarrow m_1 R \hat{i} + m_2 \left(-\frac{R}{2} \hat{i} + \frac{\sqrt{3}}{2} R \hat{j} \right) + m_3 (-R \hat{i}) + m_4 \left(-\frac{R}{2} \hat{i} - \frac{\sqrt{3}}{2} R \hat{j} \right) &= 0 \\ \Rightarrow R \{ 2(m_1 - m_3) - (m_2 + m_4) \} \hat{i} + \sqrt{3} R \{ m_2 - m_4 \} \hat{j} &= 0 \end{aligned}$$

Taking scalar product of $\hat{i}$ and $\hat{j}$ with above equations we get

$$\left. \begin{aligned} 2(m_1 - m_3) - (m_2 + m_4) &= 0 \\ m_2 &= m_4 \end{aligned} \right\} \quad (3)$$

Also we know

$$m_1 + m_2 + m_3 + m_4 = M \quad (4)$$

The equations (3) and (4) yield

$$m_1 - m_2 - m_3 = 0 \quad (5)$$

and

$$m_1 + 2m_2 + m_3 = M \quad (6)$$

Let us introduce a non- dimensional mass parameter $\mu \in (0, 3^{-1})$ such that $m_2 = m_4 = M\mu$ then from (5) & (6) we get

$$m_1 = \frac{M(1-\mu)}{2}, m_3 = \frac{M(1-3\mu)}{2}.$$

Thus the necessary conditions for four positive point masses m_k situated at the vertices of a cyclic kite formed by an equilateral triangle $P_1P_2P_4$ and an isosceles triangle $P_2P_3P_4$ with the common base P_2P_4 are

$$\left. \begin{array}{l} (a) \quad \vec{r}_1 + \vec{r}_3 = 0 \quad \vec{r}_1 + \vec{r}_2 + \vec{r}_4 = 0 \\ (b) \quad m_2 = m_4 = M\mu, \quad m_1 = \frac{M(1-\mu)}{2}, m_3 = \frac{M(1-3\mu)}{2} \\ (c) \quad m_1 - m_2 = m_3 > 0 \Rightarrow m_1 > m_2 = m_4 \end{array} \right\} \quad (7)$$

Here 5(b) justified the work of Chavella (2007).

For converse part let us prove that the diagonals of the quadrilateral $P_1P_2P_3P_4$ are perpendicular to each other. For this we have to find magnitudes of the sides of the quadrilateral. From previous results we have

$$\left. \begin{array}{l} |\vec{r}_{12}| = |\vec{r}_{21}| = |\vec{r}_2 - \vec{r}_1| = \left| -\frac{3R}{2}\hat{i} + \frac{\sqrt{3}R}{2}\hat{j} \right| = \sqrt{3}R = |\vec{r}_{14}| = |\vec{r}_{41}| = |\vec{r}_{24}| = |\vec{r}_{42}| \\ |\vec{r}_{23}| = |\vec{r}_{32}| = |\vec{r}_{34}| = |\vec{r}_{43}| = R, \quad |\vec{r}_{13}| = |\vec{r}_{31}| = 2R \end{array} \right\} \quad (8)$$

Now in the kite $P_1P_2P_3P_4$

$$\vec{r}_{13} \cdot \vec{r}_{24} = (\vec{r}_3 - \vec{r}_1) \cdot (\vec{r}_4 - \vec{r}_2) = (-R\vec{i} - R\vec{i}) \cdot \left(-\frac{R}{2}\vec{i} - \frac{\sqrt{3}}{2}R\vec{j} + \frac{R}{2}\vec{i} - \frac{\sqrt{3}}{2}R\vec{j} \right) = 0 \quad (9)$$

Therefore the diagonals P_1P_3 & P_2P_4 are perpendicular to each other and hence the quadrilateral $P_1P_2P_3P_4$ Is a cyclic kite configuration. Thus the results (7) & (9) satisfy the necessary and sufficient conditions for a quadrilateral to be a cyclic kite. These results can satisfy all the results of previous authors also.

4. The Equations of motion of four bodies of the kite in rotating frame about z- axis:

Before to derive the equations of motion of the four bodies forming cyclic kite configuration in synodic frame, let us derive first the gravitational force between two bodies (one is spherical and the other is oblate spheroid). Let A be the oblateness parameter of any oblate body then the potential V_{ij} between a spherical body of mass m_i and an oblate body of mass m_j separated by the distance r_{ij} is given by (McCuskey -1967,pp)

$$V_{ij} = -\frac{Gm_i m_j}{r_{ij}} - \frac{GA}{2r_{ij}^3} \quad (i \neq j = 1, 2, 3, 4) \quad (10)$$

Where the oblateness parameter is given by $A = I_a - I_e$, I_a is the moment of inertia of the oblate body about the polar axis and I_e is the moment of inertia about the equatorial axis.

Let the plane of motion of the four point masses (one is oblate and the others are spherical) forming kite configuration rotate with the angular velocity $\vec{\omega} = n\hat{k}$ about the z-axis ; $\hat{r}_{ij}$ be the unit vector along the line $P_i P_j$ joining the i-th and j-th point masses, then the equation of motion of the i-th point mass relative to the other three point masses is given by

$$-m_i n^2 \vec{r}_i = \sum_{j=1}^4 \frac{\partial V_{ij}}{\partial r_{ij}} \hat{r}_{ij}, \quad (i \neq j) \quad (11)$$

$$m_i n^2 \hat{r}_i + \frac{\partial V_{i1}}{\partial r_{i1}} \hat{r}_{i1} + \frac{\partial V_{i2}}{\partial r_{i2}} \hat{r}_{i2} + \frac{\partial V_{i3}}{\partial r_{i3}} \hat{r}_{i3} + \frac{\partial V_{i4}}{\partial r_{i4}} \hat{r}_{i4} = 0. \quad (12)$$

Putting $i=1,2,3,4$ in the equation (12), we get

$$\left. \begin{aligned} m_1 n^2 \vec{r}_1 + \frac{\partial V_{12}}{\partial r_{12}} \hat{r}_{12} + \frac{\partial V_{13}}{\partial r_{13}} \hat{r}_{13} + \frac{\partial V_{14}}{\partial r_{14}} \hat{r}_{14} &= 0 \\ m_2 n^2 \vec{r}_2 + \frac{\partial V_{21}}{\partial r_{21}} \hat{r}_{21} + \frac{\partial V_{23}}{\partial r_{23}} \hat{r}_{23} + \frac{\partial V_{24}}{\partial r_{24}} \hat{r}_{24} &= 0 \\ m_3 n^2 \vec{r}_3 + \frac{\partial V_{31}}{\partial r_{31}} \hat{r}_{31} + \frac{\partial V_{32}}{\partial r_{32}} \hat{r}_{32} + \frac{\partial V_{34}}{\partial r_{34}} \hat{r}_{34} &= 0 \\ m_4 n^2 \vec{r}_4 + \frac{\partial V_{41}}{\partial r_{41}} \hat{r}_{41} + \frac{\partial V_{42}}{\partial r_{42}} \hat{r}_{42} + \frac{\partial V_{43}}{\partial r_{43}} \hat{r}_{43} &= 0 \end{aligned} \right\} \quad (13)$$

As we take the first body of mass m_1 at P_1 of the kite configuration as an oblate spheroid so

$$\left. \begin{aligned} V_{12} &= -\frac{Gm_1 m_2}{r_{12}} - \frac{GA}{2r_{12}^3} = V_{21}, \\ V_{13} &= -\frac{Gm_1 m_3}{r_{13}} - \frac{GA}{2r_{13}^3} = V_{31}, \\ V_{14} &= -\frac{Gm_1 m_4}{r_{14}} - \frac{GA}{2r_{14}^3} = V_{41}. \end{aligned} \right\} \quad (14)$$

Introducing equations (14) in equations (13) one can find

$$\left. \begin{aligned} m_1 n^2 \vec{r}_1 - G \frac{\partial}{\partial r_{12}} \left[\frac{m_1 m_2}{r_{12}} + \frac{A}{2r_{12}^3} \right] \vec{r}_{12} - G \frac{\partial}{\partial r_{13}} \left[\frac{m_1 m_3}{r_{13}} + \frac{A}{2r_{13}^3} \right] \vec{r}_{13} - G \frac{\partial}{\partial r_{14}} \left[\frac{m_1 m_4}{r_{14}} + \frac{A}{2r_{14}^3} \right] \vec{r}_{14} &= 0 \\ m_2 n^2 \vec{r}_2 - G \frac{\partial}{\partial r_{21}} \left[\frac{m_2 m_1}{r_{21}} + \frac{A}{2r_{21}^3} \right] \vec{r}_{21} - G \frac{\partial}{\partial r_{23}} \left[\frac{m_2 m_3}{r_{23}} \right] \vec{r}_{23} - G \frac{\partial}{\partial r_{24}} \left[\frac{m_2 m_4}{r_{24}} \right] \vec{r}_{24} &= 0 \\ m_3 n^2 \vec{r}_3 - G \frac{\partial}{\partial r_{31}} \left[\frac{m_3 m_1}{r_{31}} + \frac{A}{2r_{31}^3} \right] \vec{r}_{31} - G \frac{\partial}{\partial r_{32}} \left[\frac{m_3 m_2}{r_{32}} \right] \vec{r}_{32} - G \frac{\partial}{\partial r_{34}} \left[\frac{m_3 m_4}{r_{34}} \right] \vec{r}_{34} &= 0 \\ m_4 n^2 \vec{r}_4 - G \frac{\partial}{\partial r_{41}} \left[\frac{m_4 m_1}{r_{41}} + \frac{A}{2r_{41}^3} \right] \vec{r}_{41} - G \frac{\partial}{\partial r_{42}} \left[\frac{m_4 m_2}{r_{42}} \right] \vec{r}_{42} - G \frac{\partial}{\partial r_{43}} \left[\frac{m_4 m_3}{r_{43}} \right] \vec{r}_{43} &= 0 \end{aligned} \right\} \quad (15)$$

$$\begin{aligned} n^2 \vec{r}_1 + G \left[\frac{m_2}{r_{12}^2} + \frac{3A}{2m_1 r_{12}^4} \right] \frac{\vec{r}_{12}}{r_{12}} + G \left[\frac{m_3}{r_{13}^2} + \frac{3A}{2m_1 r_{13}^4} \right] \frac{\vec{r}_{13}}{r_{13}} + G \left[\frac{m_4}{r_{14}^2} + \frac{3A}{2m_1 r_{14}^4} \right] \frac{\vec{r}_{14}}{r_{14}} &= 0 \\ n^2 \vec{r}_2 + G \left[\frac{m_1}{r_{21}^2} + \frac{3A}{2m_2 r_{21}^4} \right] \frac{\vec{r}_{21}}{r_{21}} + G \left[\frac{m_3}{r_{23}^2} \right] \frac{\vec{r}_{23}}{r_{23}} + G \left[\frac{m_4}{r_{24}^2} \right] \frac{\vec{r}_{24}}{r_{24}} &= 0 \\ n^2 \vec{r}_3 + G \left[\frac{m_1}{r_{31}^2} + \frac{3A}{2m_3 r_{31}^4} \right] \frac{\vec{r}_{31}}{r_{31}} + G \left[\frac{m_2}{r_{32}^2} \right] \frac{\vec{r}_{31}}{r_{31}} + G \left[\frac{m_4}{r_{34}^2} \right] \frac{\vec{r}_{31}}{r_{31}} &= 0 \\ n^2 \vec{r}_4 + G \left[\frac{m_1}{r_{41}^2} + \frac{3A}{2m_4 r_{41}^4} \right] \frac{\vec{r}_{41}}{r_{41}} + G \left[\frac{m_2}{r_{42}^2} \right] \frac{\vec{r}_{42}}{r_{42}} + G \left[\frac{m_3}{r_{43}^2} \right] \frac{\vec{r}_{43}}{r_{43}} &= 0 \end{aligned}$$

$$\begin{aligned} n^2 \vec{r}_1 + G \left[\left\{ \frac{m_2}{r_{12}^3} + \frac{3A}{2m_1 r_{12}^5} \right\} \vec{r}_{12} + \left\{ \frac{m_3}{r_{13}^3} + \frac{3A}{2m_1 r_{13}^5} \right\} \vec{r}_{13} + \left\{ \frac{m_4}{r_{14}^3} + \frac{3A}{2m_1 r_{14}^5} \right\} \vec{r}_{14} \right] &= 0 \\ n^2 \vec{r}_2 + G \left[\left\{ \frac{m_1}{r_{21}^3} + \frac{3A}{2m_2 r_{21}^5} \right\} \vec{r}_{21} + \left\{ \frac{m_3}{r_{23}^3} \right\} \vec{r}_{23} + \left\{ \frac{m_4}{r_{24}^3} \right\} \vec{r}_{24} \right] &= 0 \\ n^2 \vec{r}_3 + G \left[\left\{ \frac{m_1}{r_{31}^3} + \frac{3A}{2m_3 r_{31}^5} \right\} \vec{r}_{31} + \left\{ \frac{m_2}{r_{32}^3} \right\} \vec{r}_{32} + \left\{ \frac{m_4}{r_{34}^3} \right\} \vec{r}_{34} \right] &= 0 \\ n^2 \vec{r}_4 + G \left[\left\{ \frac{m_1}{r_{41}^3} + \frac{3A}{2m_4 r_{41}^5} \right\} \vec{r}_{41} + \left\{ \frac{m_2}{r_{42}^3} \right\} \vec{r}_{42} + \left\{ \frac{m_3}{r_{43}^3} \right\} \vec{r}_{43} \right] &= 0 \end{aligned}$$

Putting the values of m_1, m_2, m_3, m_4 from (5) we get,

$$\begin{aligned}
 n^2 \vec{r}_1 + G \left[\left\{ \frac{M\mu}{3\sqrt{3}R^3} + \frac{A}{3\sqrt{3}M(1-\mu)R^5} \right\} \vec{r}_{12} + \left\{ \frac{M(1-3\mu)}{16R^3} + \frac{3A}{32M(1-\mu)R^5} \right\} \vec{r}_{13} + \left\{ \frac{M\mu}{3\sqrt{3}R^3} + \frac{A}{3\sqrt{3}M(1-\mu)R^5} \right\} \vec{r}_{14} \right] &= \hat{0} \\
 n^2 \vec{r}_2 + G \left[\left\{ \frac{M(1-\mu)}{6\sqrt{3}R^3} + \frac{A}{6\sqrt{3}M\mu R^5} \right\} \vec{r}_{21} + \left\{ \frac{M(1-3\mu)}{2R^3} \right\} \vec{r}_{23} + \left\{ \frac{M\mu}{3\sqrt{3}R^3} \right\} \vec{r}_{24} \right] &= \hat{0} \\
 n^2 \vec{r}_3 + G \left[\left\{ \frac{M(1-\mu)}{16R^3} + \frac{3A}{32M(1-3\mu)R^5} \right\} \vec{r}_{31} + \left\{ \frac{M\mu}{R^3} \right\} \vec{r}_{32} + \left\{ \frac{M\mu}{R^3} \right\} \vec{r}_{34} \right] &= \hat{0} \\
 n^2 \vec{r}_4 + G \left[\left\{ \frac{M(1-\mu)}{6\sqrt{3}R^3} + \frac{A}{6\sqrt{3}M\mu R^5} \right\} \vec{r}_{41} + \left\{ \frac{M\mu}{3\sqrt{3}R^3} \right\} \vec{r}_{42} + \left\{ \frac{M(1-3\mu)}{2R^3} \right\} \vec{r}_{43} \right] &= \hat{0}
 \end{aligned}$$

Now choosing units of masses lengths and forces in such a way that

$$m_1 + m_2 + m_3 + m_4 = M = 1, 2R = 1, G = 1,$$

then the above equations reduced to

$$\begin{aligned}
 n^2 \vec{r}_1 + \left[\left\{ \frac{8\mu}{3\sqrt{3}} + \frac{32A}{3\sqrt{3}(1-\mu)} \right\} \vec{r}_{12} + \left\{ \frac{(1-3\mu)}{2} + \frac{3A}{(1-\mu)} \right\} \vec{r}_{13} + \left\{ \frac{8\mu}{3\sqrt{3}} + \frac{32A}{9\sqrt{3}(1-\mu)} \right\} \vec{r}_{14} \right] &= \hat{0} \\
 n^2 \vec{r}_2 + \left[\left\{ \frac{4(1-\mu)}{3\sqrt{3}} + \frac{16A}{3\sqrt{3}\mu} \right\} \vec{r}_{21} + \{4(1-3\mu)\} \vec{r}_{23} + \left\{ \frac{8\mu}{3\sqrt{3}} \right\} \vec{r}_{24} \right] &= \hat{0} \\
 n^2 \vec{r}_3 + \left[\left\{ \frac{(1-\mu)}{2} + \frac{3A}{(1-3\mu)} \right\} \vec{r}_{31} + \{8\mu\} \vec{r}_{32} + \{8\mu\} \vec{r}_{34} \right] &= \hat{0} \\
 n^2 \vec{r}_4 + \left[\left\{ \frac{4(1-\mu)}{3\sqrt{3}} + \frac{16A}{3\sqrt{3}\mu} \right\} \vec{r}_{41} + \left\{ \frac{8\mu}{3\sqrt{3}} \right\} \vec{r}_{42} + \{4(1-3\mu)\} \vec{r}_{43} \right] &= \hat{0}
 \end{aligned}$$

$\Rightarrow$

$$\begin{aligned}
 n^2 \vec{r}_1 + \left[\left\{ \frac{8\mu}{3\sqrt{3}} + \frac{32A}{3\sqrt{3}(1-\mu)} \right\} (\vec{r}_2 - \vec{r}_1) + \left\{ \frac{(1-3\mu)}{2} + \frac{3A}{(1-\mu)} \right\} (\vec{r}_3 - \vec{r}_1) + \left\{ \frac{8\mu}{3\sqrt{3}} + \frac{32A}{9\sqrt{3}(1-\mu)} \right\} (\vec{r}_4 - \vec{r}_1) \right] &= \hat{0} \\
 n^2 \vec{r}_2 + \left[\left\{ \frac{4(1-\mu)}{3\sqrt{3}} + \frac{16A}{3\sqrt{3}\mu} \right\} (\vec{r}_1 - \vec{r}_2) + \{4(1-3\mu)\} (\vec{r}_3 - \vec{r}_2) + \left\{ \frac{8\mu}{3\sqrt{3}} \right\} (\vec{r}_4 - \vec{r}_2) \right] &= \hat{0} \\
 n^2 \vec{r}_3 + \left[\left\{ \frac{(1-\mu)}{2} + \frac{3A}{(1-3\mu)} \right\} (\vec{r}_1 - \vec{r}_3) + \{8\mu\} (\vec{r}_2 - \vec{r}_3) + \{8\mu\} (\vec{r}_4 - \vec{r}_3) \right] &= \hat{0} \\
 n^2 \vec{r}_4 + \left[\left\{ \frac{4(1-\mu)}{3\sqrt{3}} + \frac{16A}{3\sqrt{3}\mu} \right\} (\vec{r}_1 - \vec{r}_4) + \left\{ \frac{8\mu}{3\sqrt{3}} \right\} (\vec{r}_2 - \vec{r}_4) + \{4(1-3\mu)\} (\vec{r}_3 - \vec{r}_4) \right] &= \hat{0}
 \end{aligned}$$

Now by arranging the above four equations in the form of linear combinations of $\vec{r}_1, \vec{r}_2, \vec{r}_3, \vec{r}_4$ one can find the equations of motion of four bodies forming kite as

$$\left. \begin{aligned} & \left[n^2 - \frac{2}{3\sqrt{3}}(8\mu + \frac{32A}{1-\mu}) - (\frac{1-3\mu}{2} + \frac{3A}{1-\mu}) \right] \vec{r}_1 + \frac{1}{3\sqrt{3}}(8\mu + \frac{32A}{1-\mu}) \vec{r}_2 + (\frac{1-3\mu}{2} + \frac{3A}{1-\mu}) \vec{r}_3 + \frac{1}{3\sqrt{3}}(8\mu + \frac{32A}{1-\mu}) \vec{r}_4 = \hat{0} \\ & \frac{4}{3\sqrt{3}}(1-\mu + \frac{4A}{\mu}) \vec{r}_1 + \left[n^2 - \frac{4}{3\sqrt{3}}(1-\mu + \frac{4A}{\mu}) - 4(1-3\mu) - \frac{8\mu}{3\sqrt{3}} \right] \vec{r}_2 + 4(1-3\mu) \vec{r}_3 + \frac{8\mu}{3\sqrt{3}} \vec{r}_4 = \hat{0} \\ & (\frac{1-\mu}{2} + \frac{3A}{1-3\mu}) \vec{r}_1 + 8\mu \vec{r}_2 + \left[n^2 - (\frac{1-\mu}{2} + \frac{3A}{1-3\mu}) - 16\mu \right] \vec{r}_3 + 8\mu \vec{r}_4 = \hat{0} \\ & \frac{4}{3\sqrt{3}}(1-\mu + \frac{4A}{\mu}) \vec{r}_1 + 8\mu \vec{r}_2 + 4(1-3\mu) \vec{r}_3 + \left[n^2 - \frac{4}{3\sqrt{3}}(1-3\mu + \frac{4A}{\mu}) - 4(1-3\mu) \right] \vec{r}_4 = 0 \end{aligned} \right\} (16)$$

5. Mean Motion of the rotating frame:

To find the mean motion 'n' of the rotating frame we have to express n in terms of other parameters. In the above four equations of (16) there are three parameters n, μ, A other than four position vectors $\vec{r}_1, \vec{r}_2, \vec{r}_3, \vec{r}_4$.

Thus by eliminating $\vec{r}_1, \vec{r}_2, \vec{r}_3, \vec{r}_4$ from the above four equations we get

$$\begin{vmatrix} a & \frac{8}{3\sqrt{3}}(\mu + \frac{4A}{1-\mu}) & (\frac{1-3\mu}{2} + \frac{3A}{1-\mu}) & \frac{8}{3\sqrt{3}}(\mu + \frac{4A}{1-\mu}) \\ \frac{4}{3\sqrt{3}}(1-\mu + \frac{4A}{\mu}) & b & 4(1-3\mu) & \frac{8\mu}{3\sqrt{3}} \\ \frac{1-\mu}{2} + \frac{3A}{1-3\mu} & 8\mu & c & 8\mu \\ \frac{4}{3\sqrt{3}}(1-\mu + \frac{4A}{\mu}) & 8\mu & 4(1-3\mu) & b \end{vmatrix} = 0 \quad (17)$$

Where

$$a = n^2 - \frac{16}{3\sqrt{3}}(\mu + \frac{4A}{1-\mu}) - (\frac{1-3\mu}{2} + \frac{3A}{1-\mu}), \quad b = n^2 - \frac{4}{3\sqrt{3}}(1-\mu + \frac{4A}{\mu}) - 4(1-3\mu), \quad c = n^2 - (\frac{1-\mu}{2} + \frac{3A}{1-3\mu}) - 16\mu.$$

The equation (16) gives the mean motion 'n' of the rotating (synodic) frame as a function of μ and A . i.e.

$$n = f(\mu, A).$$

6: Exact domain of the mass parameter μ in classical case (for $A=0$):

If we put $A=0$, $\mu = 1/3$ then $m_1 = m_2 = m_4 = 1/3$ & $m_3 = 0$. Thus the cyclic kite configuration converges to the cyclic equilateral triangular configuration and from (16) we have

$$\begin{vmatrix} n^2 - \frac{16}{9\sqrt{3}} & \frac{8}{9\sqrt{3}} & 0 & \frac{8}{9\sqrt{3}} \\ \frac{8}{9\sqrt{3}} & n^2 - \frac{64}{9\sqrt{3}} & 0 & \frac{8}{9\sqrt{3}} \\ \frac{1}{3} & \frac{8}{3} & n^2 - \frac{17}{3} & \frac{8}{3} \\ \frac{8}{9\sqrt{3}} & \frac{8}{3} & 0 & n^2 - \frac{64}{9\sqrt{3}} \end{vmatrix} = 0 \quad (17)$$

$$\Rightarrow n = 0.8726844, 1.79568456, 2.29162458, 2.33047614.$$

But equilateral triangular configuration rotates with unit mean motion so these values of mean motion are invalid and consequently the kite configuration does not exist for $\mu = 1/3$. Thus the exact domain of μ is $(0, 3^{-1})$.

7. The data produced to analyze the effect of oblateness on the mean motion of the rotating frame by using Python.

S/No.	μ	A	n	S/No.	μ	A	n	S/No.	μ	A	n
1	0.01	0.01	1.094904	18	0.01	0.18	1.530321	35	0.02	0.15	1.544661
2	0.01	0.02	1.141577	19	0.01	0.19	1.550052	36	0.02	0.16	1.565094
3	0.01	0.03	1.177312	20	0.01	0.2	1.569512	37	0.02	0.17	1.585149
4	0.01	0.04	1.208292	21	0.02	0.01	1.120484	38	0.02	0.18	1.604858
5	0.01	0.05	1.236655	22	0.02	0.02	1.183872	39	0.02	0.19	1.624246
6	0.01	0.06	1.263339	23	0.02	0.03	1.230685	40	0.02	0.2	1.643336
7	0.01	0.07	1.288829	24	0.02	0.04	1.268953	41	0.03	0.01	1.137155
8	0.01	0.08	1.313402	25	0.02	0.05	1.302219	42	0.03	0.02	1.212038
9	0.01	0.09	1.337233	26	0.02	0.06	1.332265	43	0.03	0.03	1.268233
10	0.01	0.1	1.360438	27	0.02	0.07	1.360082	44	0.03	0.04	1.313568
11	0.01	0.11	1.3831	28	0.02	0.08	1.386267	45	0.03	0.05	1.352094
12	0.01	0.12	1.40528	29	0.02	0.09	1.411199	46	0.03	0.06	1.386074
13	0.01	0.13	1.427024	30	0.02	0.1	1.435133	47	0.03	0.07	1.416857
14	0.01	0.14	1.448371	31	0.02	0.11	1.458247	48	0.03	0.08	1.445293
15	0.01	0.15	1.46935	32	0.02	0.12	1.480672	49	0.03	0.09	1.471941
16	0.01	0.16	1.489988	33	0.02	0.13	1.502502	50	0.03	0.1	1.497184
17	0.01	0.17	1.510305	34	0.02	0.14	1.523813	51	0.03	0.11	1.521293

S/No.	μ	A	n	S/No.	μ	A	n	S/No.	μ	A	n
52	0.03	0.12	1.544466	90	0.05	0.1	1.598497	127	0.07	0.07	1.575389
53	0.03	0.13	1.566852	91	0.05	0.11	1.625267	128	0.07	0.08	1.61423
54	0.03	0.14	1.588563	92	0.05	0.12	1.650594	129	0.07	0.09	1.649398
55	0.03	0.15	1.609687	93	0.05	0.13	1.674722	130	0.07	0.1	1.681628
56	0.03	0.16	1.630295	94	0.05	0.14	1.697839	131	0.07	0.11	1.711477
57	0.03	0.17	1.650442	95	0.05	0.15	1.720093	132	0.07	0.12	1.739373
58	0.03	0.18	1.670176	96	0.05	0.16	1.741602	133	0.07	0.13	1.765648
59	0.03	0.19	1.689534	97	0.05	0.17	1.762461	134	0.07	0.14	1.790563
60	0.03	0.2	1.708548	98	0.05	0.18	1.782747	135	0.07	0.15	1.814324
61	0.04	0.01	1.150072	99	0.05	0.19	1.802523	136	0.07	0.16	1.837096
62	0.04	0.02	1.233274	100	0.05	0.2	1.821842	137	0.07	0.17	1.859013
63	0.04	0.03	1.297213	101	0.06	0.01	1.170298	138	0.07	0.18	1.880183
64	0.04	0.04	1.348933	102	0.06	0.02	1.264858	139	0.07	0.19	1.900698
65	0.04	0.05	1.392536	103	0.06	0.03	1.340777	140	0.07	0.2	1.920631
66	0.04	0.06	1.430515	104	0.06	0.04	1.40338	141	0.08	0.01	1.186299
67	0.04	0.07	1.464449	105	0.06	0.05	1.456288	142	0.08	0.02	1.288356
68	0.04	0.08	1.495376	106	0.06	0.06	1.502015	143	0.08	0.03	1.373163
69	0.04	0.09	1.524002	107	0.06	0.07	1.54233	144	0.08	0.04	1.444697
70	0.04	0.1	1.55082	108	0.06	0.08	1.578491	145	0.08	0.05	1.505885
71	0.04	0.11	1.576187	109	0.06	0.09	1.61141	146	0.08	0.06	1.558951
72	0.04	0.12	1.600364	110	0.06	0.1	1.641756	147	0.08	0.07	1.605604
73	0.04	0.13	1.623549	111	0.06	0.11	1.670028	148	0.08	0.08	1.647163
74	0.04	0.14	1.645894	112	0.06	0.12	1.696601	149	0.08	0.09	1.684642
75	0.04	0.15	1.667516	113	0.06	0.13	1.721766	150	0.08	0.1	1.718827
76	0.04	0.16	1.68851	114	0.06	0.14	1.745747	151	0.08	0.11	1.750326
77	0.04	0.17	1.708952	115	0.06	0.15	1.768722	152	0.08	0.12	1.779614
78	0.04	0.18	1.728902	116	0.06	0.16	1.790831	153	0.08	0.13	1.807065
79	0.04	0.19	1.748412	117	0.06	0.17	1.812189	154	0.08	0.14	1.832974
80	0.04	0.2	1.767524	117	0.06	0.17	1.812189	155	0.08	0.15	1.857577
81	0.05	0.01	1.160882	118	0.06	0.18	1.832889	156	0.08	0.16	1.881063
82	0.05	0.02	1.250422	119	0.06	0.19	1.853008	157	0.08	0.17	1.903586
83	0.05	0.03	1.320836	120	0.06	0.2	1.872608	158	0.08	0.18	1.925272
84	0.05	0.04	1.378281	121	0.07	0.01	1.178694	159	0.08	0.19	1.946224
85	0.05	0.05	1.426677	122	0.07	0.02	1.277346	160	0.08	0.2	1.966529
86	0.05	0.06	1.468583	123	0.07	0.03	1.358013	161	0.09	0.01	1.193269
87	0.05	0.07	1.505708	124	0.07	0.04	1.42529	162	0.09	0.02	1.298203
88	0.05	0.08	1.539227	125	0.07	0.05	1.482454	163	0.09	0.03	1.386649
89	0.05	0.09	1.56996	126	0.07	0.06	1.531899	164	0.09	0.04	1.462074

S/No.	μ	A	n	S/No.	μ	A	n	S/No.	μ	A	n
165	0.09	0.05	1.527069	203	0.11	0.03	1.409781	241	0.13	0.01	1.216902
166	0.09	0.06	1.583665	204	0.11	0.04	1.492027	242	0.13	0.02	1.329853
167	0.09	0.07	1.633476	205	0.11	0.05	1.564031	243	0.13	0.03	1.429154
168	0.09	0.08	1.677795	206	0.11	0.06	1.627409	244	0.13	0.04	1.517101
169	0.09	0.09	1.717653	207	0.11	0.07	1.683516	245	0.13	0.05	1.595319
170	0.09	0.1	1.753869	208	0.11	0.08	1.733504	246	0.13	0.06	1.66507
171	0.09	0.11	1.787095	209	0.11	0.09	1.778356	247	0.13	0.07	1.727406
172	0.09	0.12	1.817849	210	0.11	0.1	1.818909	248	0.13	0.08	1.783257
173	0.09	0.13	1.846541	211	0.11	0.11	1.855862	249	0.13	0.09	1.833466
174	0.09	0.14	1.873504	212	0.11	0.12	1.889798	250	0.13	0.1	1.878796
175	0.09	0.15	1.899001	213	0.11	0.13	1.921201	251	0.13	0.11	1.919934
176	0.09	0.16	1.923247	214	0.11	0.14	1.950468	252	0.13	0.12	1.957491
177	0.09	0.17	1.946418	215	0.11	0.15	1.977923	253	0.13	0.13	1.991995
178	0.09	0.18	1.968656	216	0.11	0.16	2.003836	254	0.13	0.14	2.023904
179	0.09	0.19	1.990081	217	0.11	0.17	2.028428	255	0.13	0.15	2.053604
180	0.09	0.2	2.01079	218	0.11	0.18	2.051879	256	0.13	0.16	2.081421
181	0.1	0.01	1.199724	219	0.11	0.19	2.074341	257	0.13	0.17	2.107628
182	0.1	0.02	1.307116	220	0.11	0.2	2.09594	258	0.13	0.18	2.132452
183	0.1	0.03	1.398777	221	0.12	0.01	1.211462	259	0.13	0.19	2.156084
184	0.1	0.04	1.477763	222	0.12	0.02	1.322809	260	0.13	0.2	2.178682
185	0.1	0.05	1.546362	223	0.12	0.03	1.419853	261	0.14	0.01	1.222149
186	0.1	0.06	1.606396	224	0.12	0.04	1.505079	262	0.14	0.02	1.336508
187	0.1	0.07	1.659359	225	0.12	0.05	1.580289	263	0.14	0.03	1.437824
188	0.1	0.08	1.706487	226	0.12	0.06	1.646911	264	0.14	0.04	1.528252
189	0.1	0.09	1.7488	227	0.12	0.07	1.706146	265	0.14	0.05	1.609282
190	0.1	0.1	1.787135	228	0.12	0.08	1.759043	266	0.14	0.06	1.682035
191	0.1	0.11	1.822175	229	0.12	0.09	1.806527	267	0.14	0.07	1.747428
192	0.1	0.12	1.854473	230	0.12	0.1	1.849407	268	0.14	0.08	1.806265
193	0.1	0.13	1.88448	231	0.12	0.11	1.888388	269	0.14	0.09	1.859291
194	0.1	0.14	1.912558	232	0.12	0.12	1.924071	270	0.14	0.1	1.9072
195	0.1	0.15	1.939004	233	0.12	0.13	1.956968	271	0.14	0.11	1.950643
196	0.1	0.16	1.964058	234	0.12	0.14	1.987505	272	0.14	0.12	1.99022
197	0.1	0.17	1.987918	235	0.12	0.15	2.016041	273	0.14	0.13	2.026469
198	0.1	0.18	2.010745	236	0.12	0.16	2.042871	274	0.14	0.14	2.05987
199	0.1	0.19	2.032673	237	0.12	0.17	2.068242	275	0.14	0.15	2.090835
200	0.1	0.2	2.053813	238	0.12	0.18	2.092358	276	0.14	0.16	2.119722
201	0.11	0.01	1.205761	239	0.12	0.19	2.115386	277	0.14	0.17	2.146831
202	0.11	0.02	1.315271	240	0.12	0.2	2.13747	278	0.14	0.18	2.172418

S/No.	μ	A	n	S/No.	μ	A	n	S/No.	μ	A	n
279	0.14	0.19	2.196695	317	0.16	0.17	2.225662	355	0.18	0.15	2.24046
280	0.14	0.2	2.219841	318	0.16	0.18	2.252987	356	0.18	0.16	2.274934
281	0.15	0.01	1.227266	319	0.16	0.19	2.27872	357	0.18	0.17	2.306738
282	0.15	0.02	1.342872	320	0.16	0.2	2.303089	358	0.18	0.18	2.336247
283	0.15	0.03	1.445993	321	0.17	0.01	1.237342	359	0.18	0.19	2.363791
284	0.15	0.04	1.538682	322	0.17	0.02	1.35508	360	0.18	0.2	2.389659
285	0.15	0.05	1.622331	323	0.17	0.03	1.461305	361	0.19	0.01	1.247584
286	0.15	0.06	1.69795	324	0.17	0.04	1.557954	362	0.19	0.02	1.367152
287	0.15	0.07	1.766335	325	0.17	0.05	1.646303	363	0.19	0.03	1.476024
288	0.15	0.08	1.828174	326	0.17	0.06	1.727226	364	0.19	0.04	1.576089
289	0.15	0.09	1.884099	327	0.17	0.07	1.80135	365	0.19	0.05	1.668595
290	0.15	0.1	1.93472	328	0.17	0.08	1.869158	366	0.19	0.06	1.754369
291	0.15	0.11	1.980628	329	0.17	0.09	1.931068	367	0.19	0.07	1.833953
292	0.15	0.12	2.022393	330	0.17	0.1	1.987483	368	0.19	0.08	1.907692
293	0.15	0.13	2.060548	331	0.17	0.11	2.038821	369	0.19	0.09	1.975809
294	0.15	0.14	2.095584	332	0.17	0.12	2.08553	370	0.19	0.1	2.038474
295	0.15	0.15	2.12794	333	0.17	0.13	2.128076	371	0.19	0.11	2.095872
296	0.15	0.16	2.158	334	0.17	0.14	2.166931	372	0.19	0.12	2.148258
297	0.15	0.17	2.186095	335	0.17	0.15	2.202555	373	0.19	0.13	2.195972
298	0.15	0.18	2.21251	336	0.17	0.16	2.235376	374	0.19	0.14	2.239432
299	0.15	0.19	2.237482	337	0.17	0.17	2.265785	375	0.19	0.15	2.279093
300	0.15	0.2	2.261215	338	0.17	0.18	2.294129	376	0.19	0.16	2.315411
301	0.16	0.01	1.232311	339	0.17	0.19	2.320707	377	0.19	0.17	2.348816
302	0.16	0.02	1.349034	340	0.17	0.2	2.345778	378	0.19	0.18	2.379699
303	0.16	0.03	1.453781	341	0.18	0.01	1.242415	379	0.19	0.19	2.408408
304	0.16	0.04	1.548535	342	0.18	0.02	1.361091	380	0.19	0.2	2.43525
305	0.16	0.05	1.634618	343	0.18	0.03	1.46868	381	0.2	0.01	1.252907
306	0.16	0.06	1.712959	344	0.18	0.04	1.567086	382	0.2	0.02	1.373348
307	0.16	0.07	1.784258	345	0.18	0.05	1.657561	383	0.2	0.03	1.483463
308	0.16	0.08	1.849093	346	0.18	0.06	1.740946	384	0.2	0.04	1.585144
309	0.16	0.09	1.907986	347	0.18	0.07	1.817815	385	0.2	0.05	1.679655
310	0.16	0.1	1.961446	348	0.18	0.08	1.888575	386	0.2	0.06	1.767842
311	0.16	0.11	2.009986	349	0.18	0.09	1.953544	387	0.2	0.07	1.850249
312	0.16	0.12	2.054125	350	0.18	0.1	2.013016	388	0.2	0.08	1.927186
313	0.16	0.13	2.094372	351	0.18	0.11	2.067306	389	0.2	0.09	1.998781
314	0.16	0.14	2.131218	352	0.18	0.12	2.116773	390	0.2	0.1	2.065037
315	0.16	0.15	2.165115	353	0.18	0.13	2.161826	391	0.2	0.11	2.125914
316	0.16	0.16	2.196476	354	0.18	0.14	2.202904	392	0.2	0.12	2.181445

S/No.	μ	A	n	S/No.	μ	A	n
393	0.2	0.13	2.231847	397	0.2	0.17	2.392346
394	0.2	0.14	2.27755	398	0.2	0.18	2.424924
395	0.2	0.15	2.319139	399	0.2	0.19	2.455271
396	0.2	0.16	2.357225	400	0.2	0.2	2.483627

8 .Graphical study of the effect of oblateness on the mean motion of the rotating frame

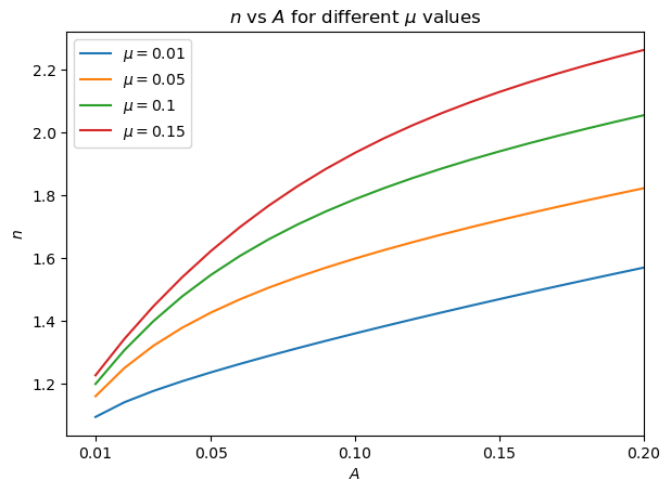


Fig-2 Graph of A vs n for different values μ

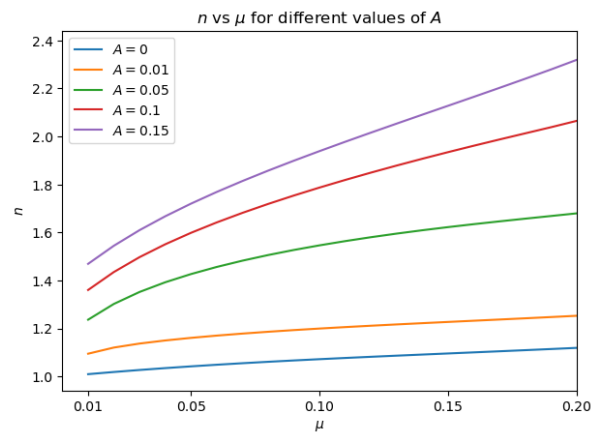


Fig-3 Graph of μ vs n for different values of A

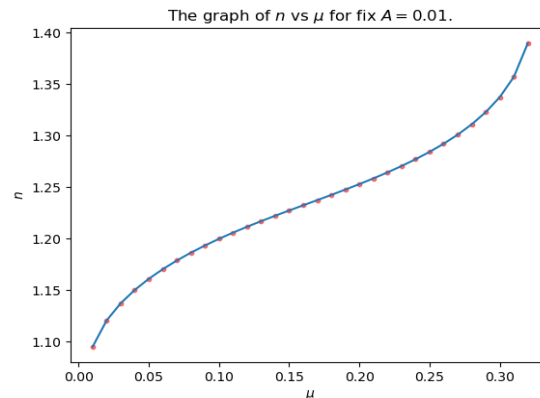


Fig-4 (Graph of μ vs n for a fixed value of A)

9. Summary

In the theorem of section 2.1 two masses situated on the opposite sides of the axis of symmetry are shown to be equal and supposed to be the mass ratio μ of either of the equal masses to the total mass $M = m_1 + m_2 + m_3 + m_4$ of kite configuration. Later on other two unequal masses lying on the axis of symmetry have been represented explicitly. Not only that, in section 2.2, it has been proved independently that for a cyclic kite configuration $P_1P_2P_3P_4$ (in figure-1) the diagonals P_1OP_3 and P_2P_4 are perpendicular to each other and at least one of them is the axis of symmetry of the kite. In our work; the diagonal P_1OP_3 is supposed to be the axis of symmetry and masses m_2 & m_4 lying on the opposite side of the axis of symmetry to be equal that is $m_2 = m_4 = M\mu$, $m_1 = M(1 - \mu)/2$, $m_3 = M(1 - 3\mu)/2$. In section 3 the equations of motion of each mass of the kite relative to the other three bodies have been derived, then by eliminating the position vectors of the four bodies, the equations for the mean motion was established in equation (15). By using Mathematica 11 the mean motion n has been expressed in terms of μ & A i.e.; $n = f(\mu, A)$. In section 3.1 we have analytically defined the exact domain of the mass parameter μ as $(0, 3^{-1})$ in classical case. In section -4 we have tabulated the data for μ, n, A for graphical analysis. In figure -2 we have shown that mean motion n increases with the increase of mass parameter μ when $A = 0$. In figures -3 & 4 keeping μ fixed, it has been shown that n increases with the increase of A but n becomes imaginary if $\mu \geq 0.2$ whereas as analytically $0 < \mu < 1/3$. When $\mu > 0.2$ is taken, the graph of n versus μ becomes discontinuous. Thus in the present paper the effect of oblateness have been discussed analytically and graphically both on the mean motion of the rotating frame. For further study of the satellite motion in the gravitational field of the kite configuration it is necessary to find n .

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